

# Structure and magnetic alignment of self-assembled rod-coil molecular aggregates

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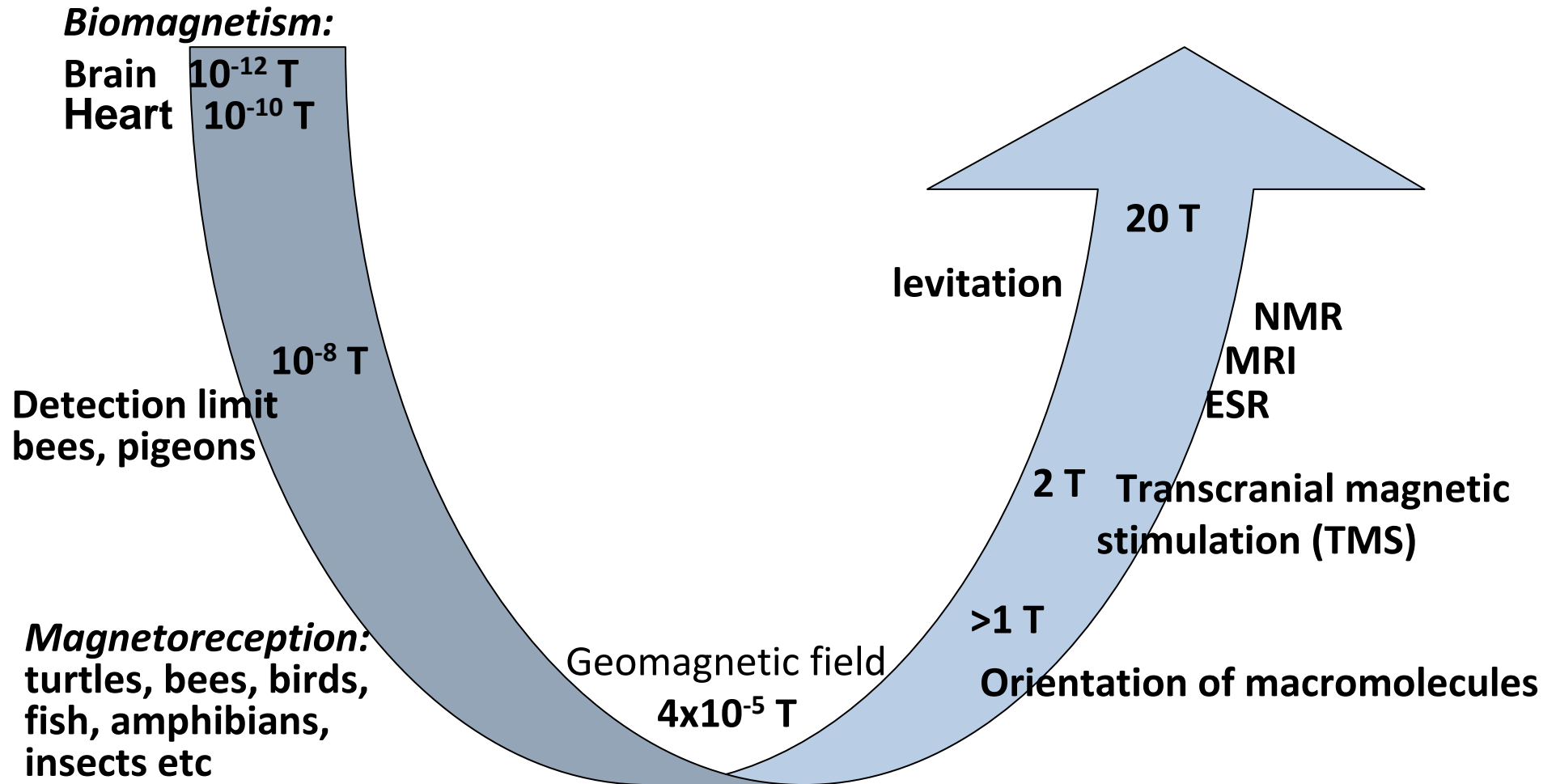
SYNEMAG - Grenoble October 17-19, 2012





or

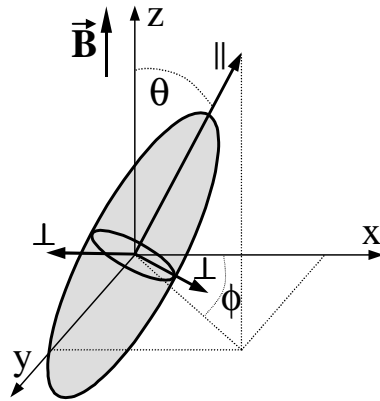
# Soft matter manipulation via magnetic field





# Magnetic Orientation

Anisotropic diamagnetic molecules in B:



$$\mathbf{m} = \frac{\chi V}{\mu_0} \cdot \mathbf{B}$$

$$\chi = \begin{pmatrix} \chi_{xx} & 0 & 0 \\ 0 & \chi_{yy} & 0 \\ 0 & 0 & \chi_{zz} \end{pmatrix}$$

$$\begin{cases} m_{\perp} = \frac{\chi_{\perp} V \cdot B}{\mu_0} & \Rightarrow & E = -\frac{\chi_{\perp} V}{2\mu_0} B^2 \\ m_{\parallel} = \frac{\chi_{\parallel} V \cdot B}{\mu_0} & \Rightarrow & E = -\frac{\chi_{\parallel} V}{2\mu_0} B^2 \end{cases}$$

$$\Delta E_{an} = -\frac{(\chi_{\parallel} - \chi_{\perp})V}{2\mu_0} B^2 = -\frac{\Delta\chi V}{2\mu_0} B^2$$



# Magnetic Field induced alignment

Requirement:

$$\frac{\Delta\chi}{2\mu_0} B^2 \geq k_B T$$

Textbook example: Benzene molecule

$$\Delta\chi = -750 \times 10^{-12} \text{ m}^3/\text{mol}$$

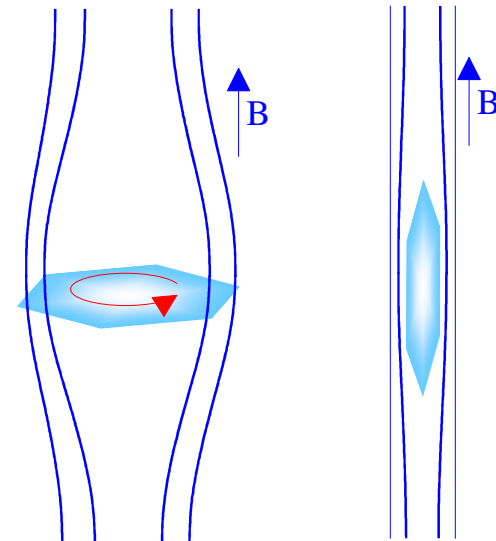
$$\Delta E_{\text{an}} = 10^{-5} k_B T \text{ (20 T, 300 K)}$$



Too small for single molecule alignment



Liquid crystals, macromolecules, polymers, molecular aggregates  
( $N|\Delta\chi|H^2 > k_B T$ )



G. Maret and K. Dransfeld (1985)



## Orientation of Non-Magnetic Materials in a Magnetic Field

- Most molecules have a diamagnetic anisotropy,  $\Delta\chi$
- For non-interacting molecules:  
 $|\Delta\chi|H^2 \ll k_B T$

**Orienting energy  $\ll$  thermal energy  
Weak orientation**

**has uses in the study of solutions:**

- flexibility of polymers
- onset of cooperativity (liquid crystals)
- solution properties (ionic solutions)

- When N molecules behave in unison their  $\Delta\chi$ s can add giving a combined anisotropy:

$$N|\Delta\chi|H^2 > k_B T$$

**Orienting energy  $>$  thermal energy  
High orientation**

$$M > 10^{10} \text{ Daltons (H=10 T)}$$

Can dramatically transform:

- information
- properties

**High orientation requires cooperativity: large, ordered assemblies such as crystals, liquid crystals and semi-rigid polymers**



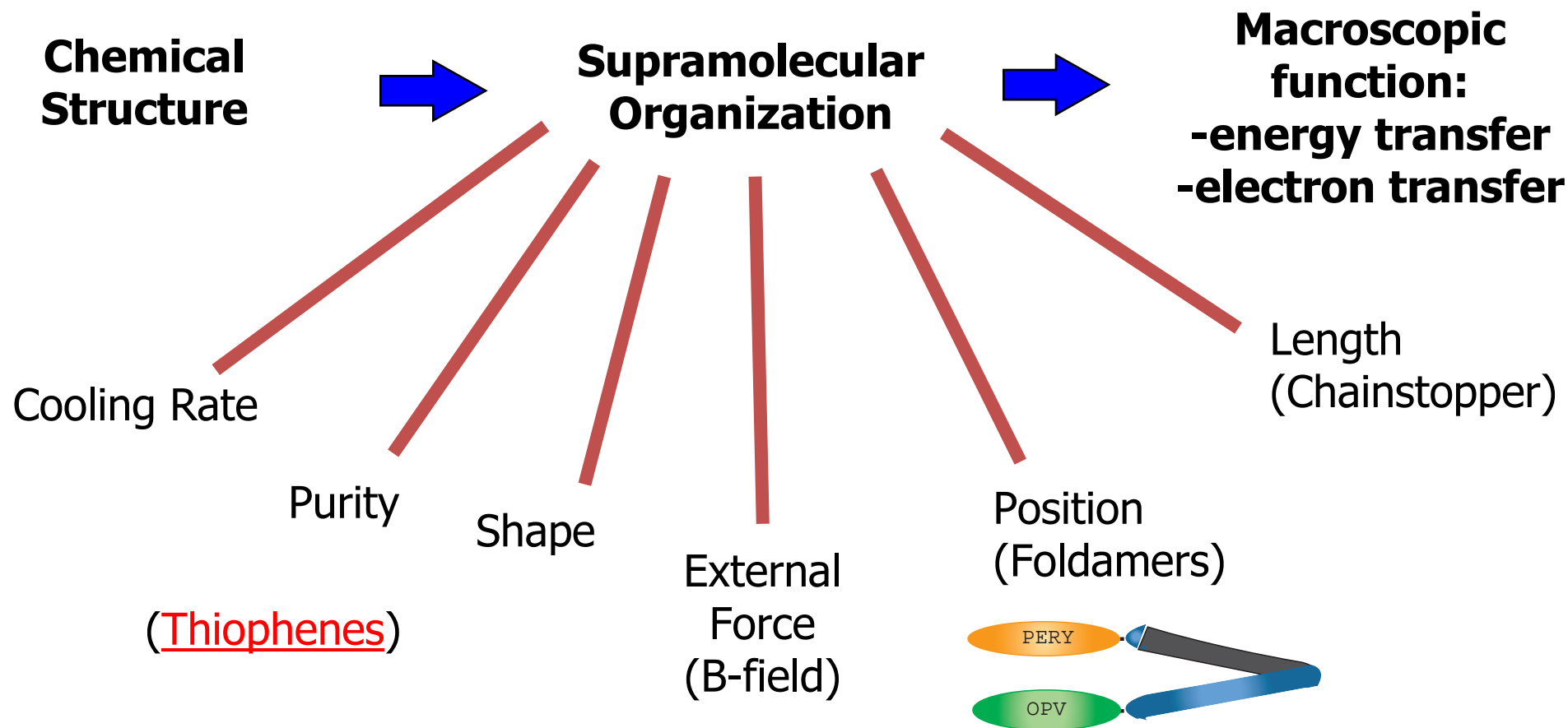
# Several materials can be aligned

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Liquid crystals	de Gennes, The physics of liquid crystals
Polymers	Kimura, Polymer Journal ('03), ...
LC Polymers	Anwer, Polymer ('91, '93), Benicewicz, Macromol. ('98), Boamfa, PRL ('03), ...
Phospholipids	Tenforde, J. Th. Bio. ('88), Ozeki, J. Phys. Chem. ('00), ...
Nucleic acids	Maret, Phys. Rev. Lett. ('75), ...
Bacteriophages	Torbet, J. Mol. Biol. ('79); Banner, Nature ('81), ...
Fibrin	Torbet, Nature ('81); ...
Collagen	Torbet, Biochem. J. ('84); Tranquillo, J. Cell Sci. ('93), ...
Tubulin	Bras, Biophys. J. ('98); ...
Carbon nanotubes	Smith, APL ('00), Hone, APL ('00), Choi, JAP ('03), Garmestani, Adv. Mat. ('03)
Protein Crystals	Wakayama, J. Cryst. Growth ('01), ...
Dye aggregates	Shklyarevskiy, J. Chem. Phys. ('02), ...

Review: G. Maret and K. Dransfeld (1985)



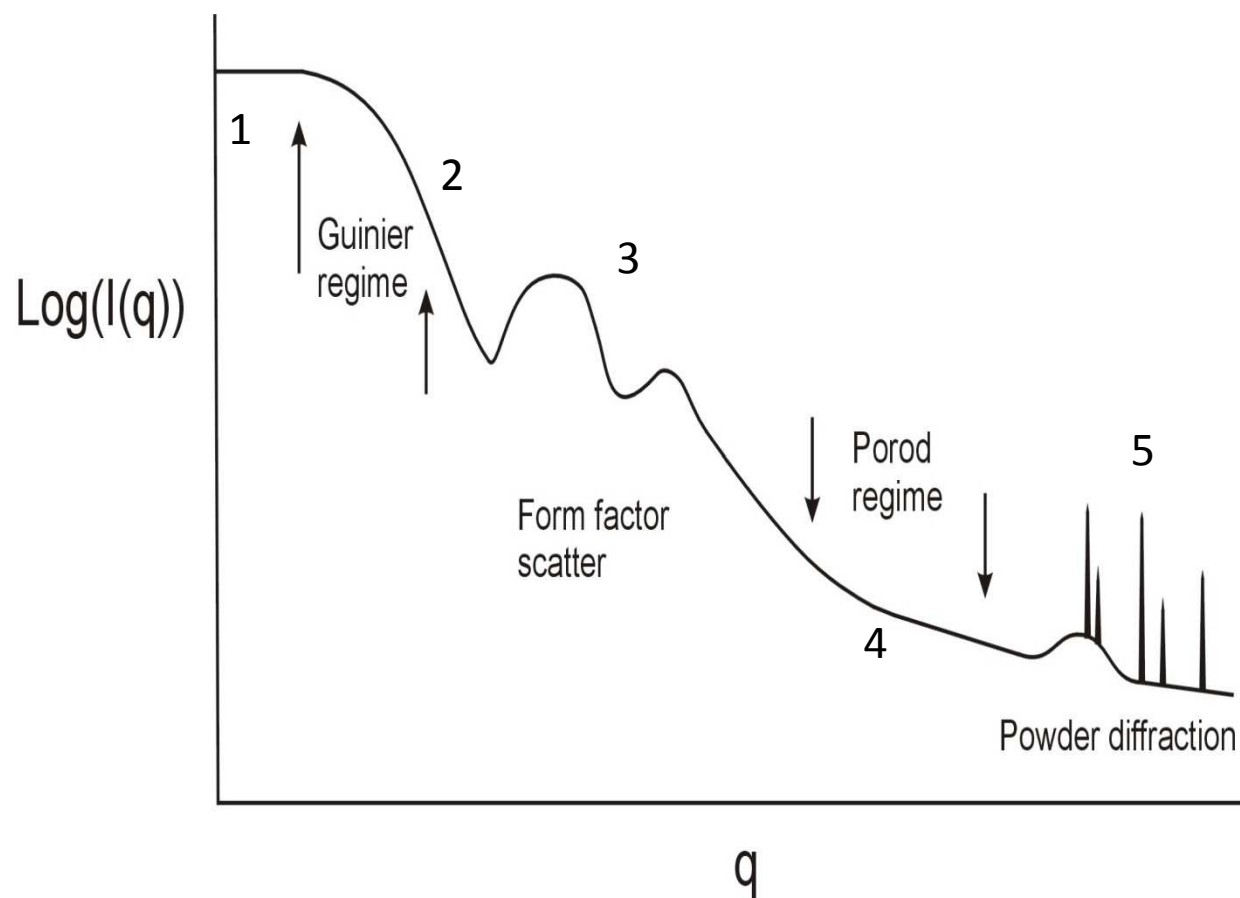


## Controlling Chemical Self-Assembly





# SAXS/WAXS



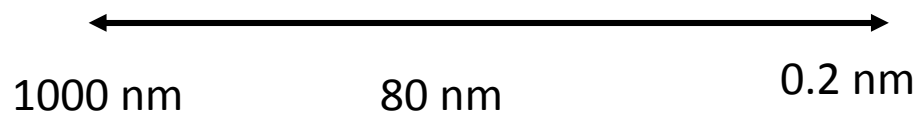
1 limit  $q \rightarrow 0$   
molecular weights

2 Guinier range  
particle size

3 particle shape  
large scale structures

4 Porod range  
particle surface

5 Intermolecular/atomic  
ordering

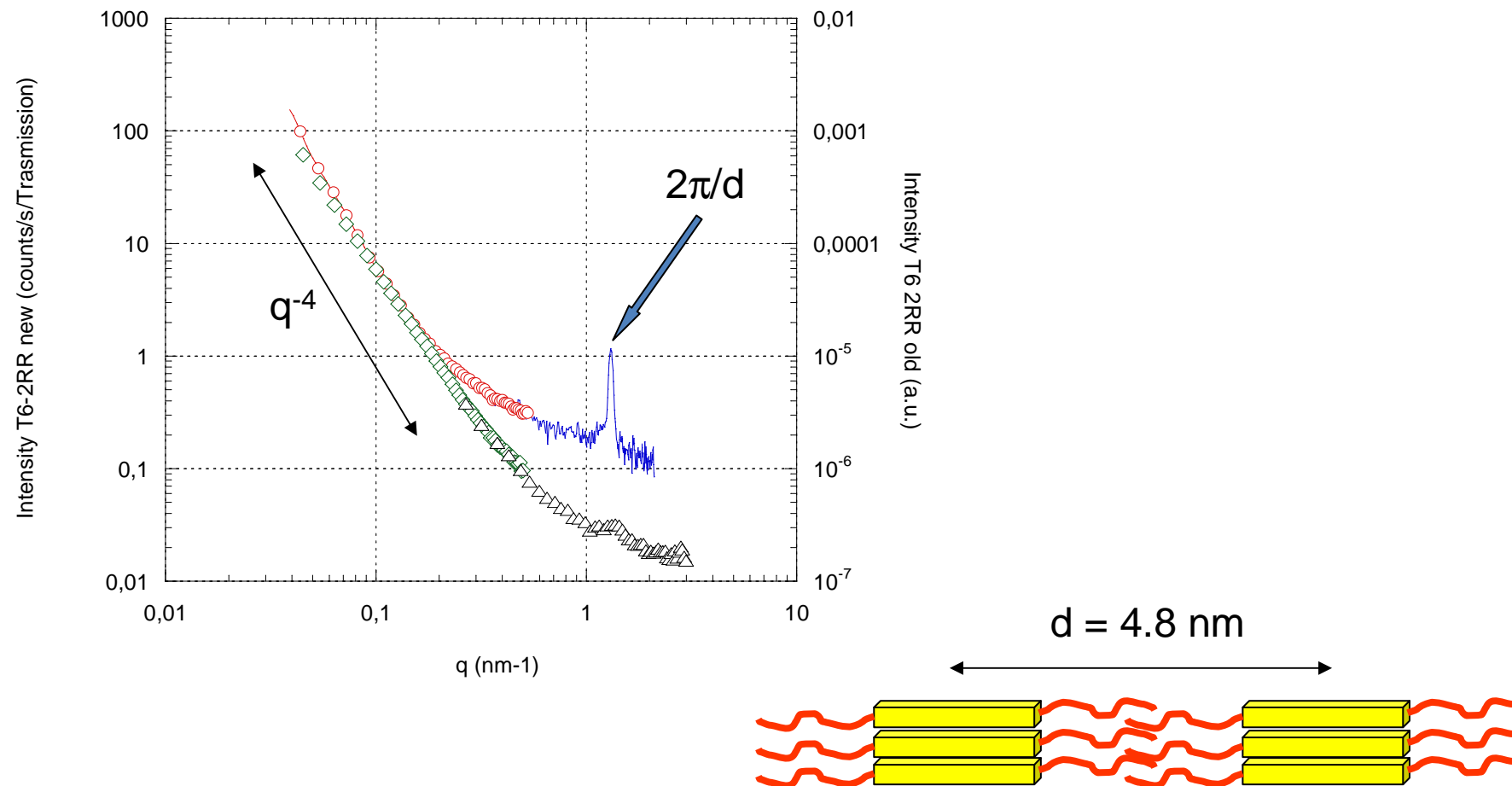






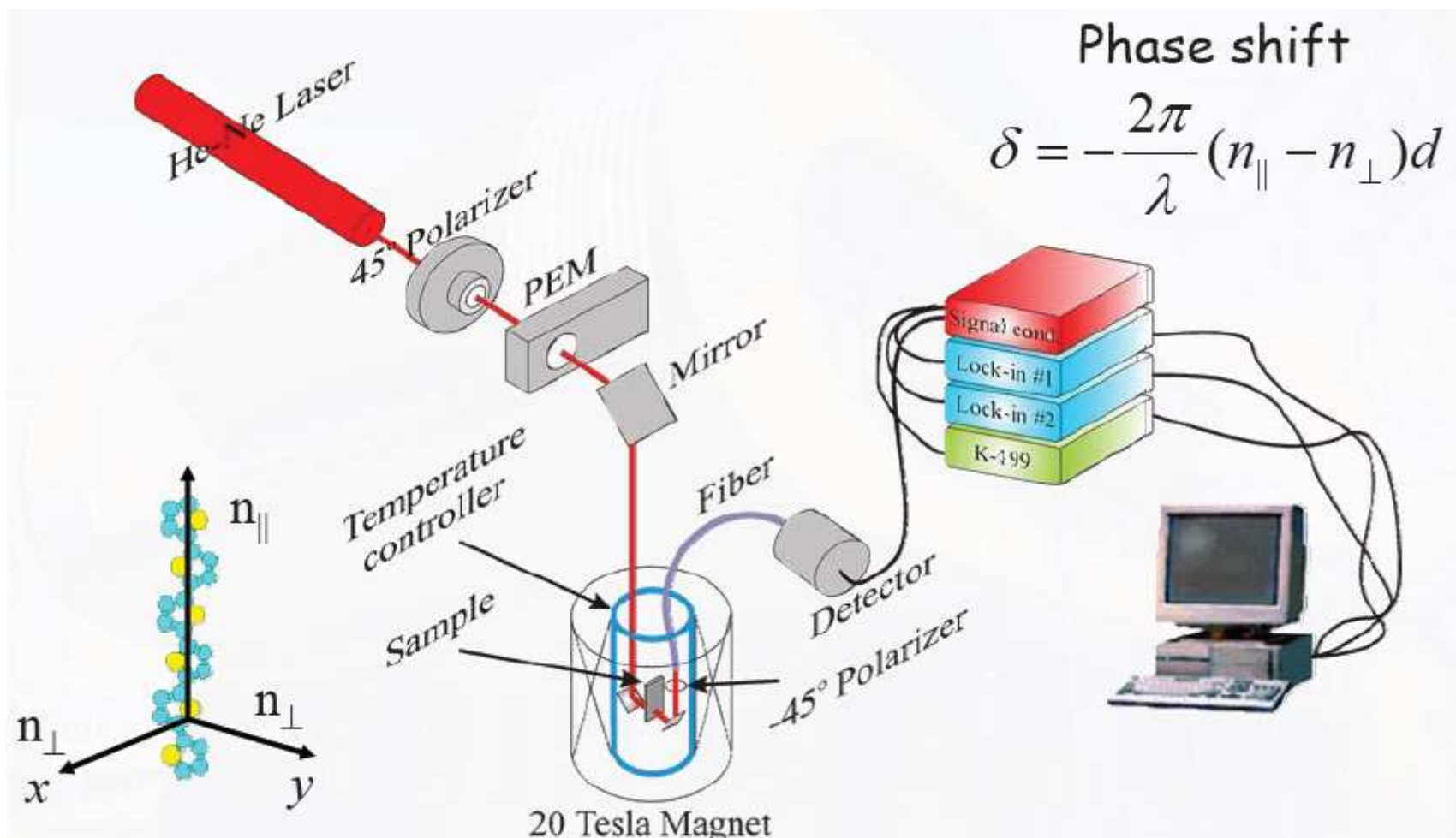


# X-ray Scattering on isotropic samples



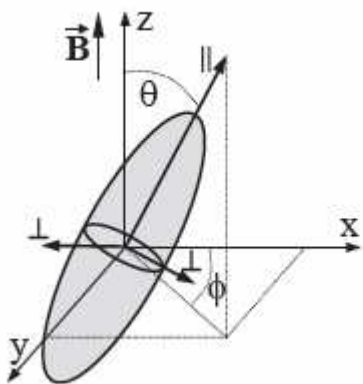
Sharp peak at 4.8 nm related to internal structure

# In-situ birefringence setup





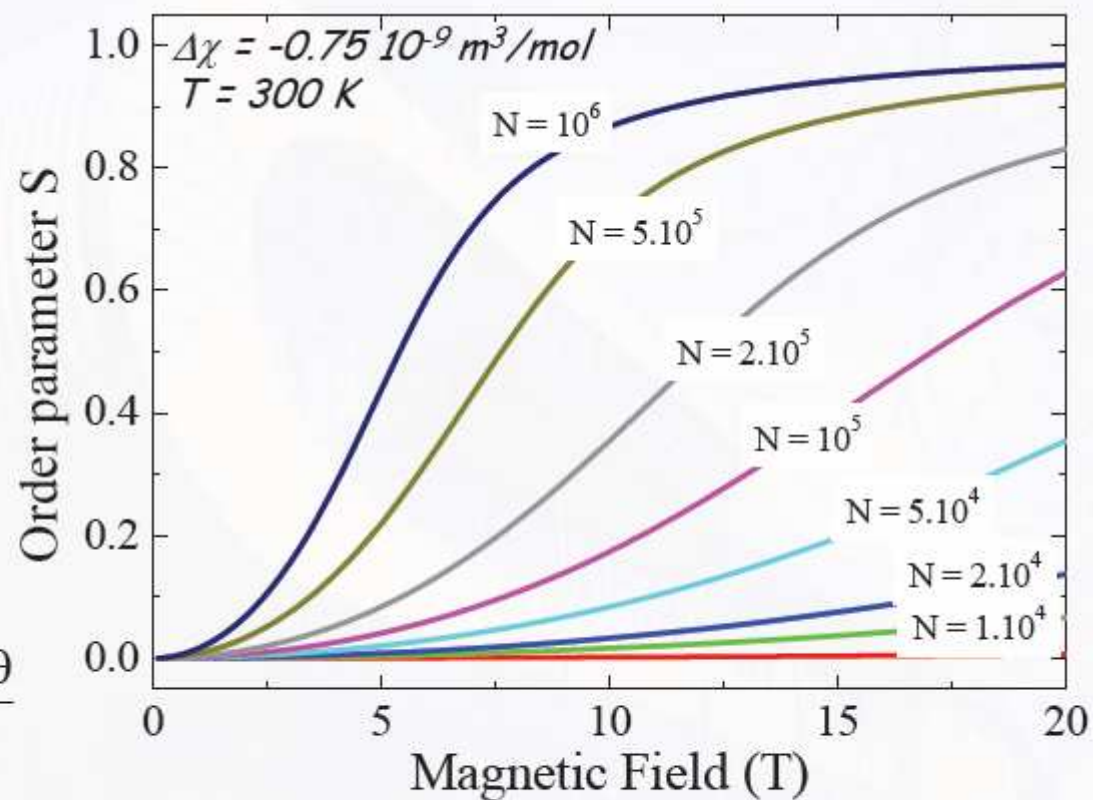
# Birefringence – degree of alignment and size dependence



$$S = \frac{3 \cdot \langle \cos^2 \theta \rangle - 1}{2}$$

$$\langle \cos^2 \theta \rangle = \frac{\int \cos^2 \theta \cdot f(\theta) \cdot \sin \theta \cdot d\theta}{\int f(\theta) \cdot \sin \theta \cdot d\theta}$$

$$f(\theta) \sim \exp \left[ \frac{N \cdot \Delta\chi \cdot B^2 \cdot \cos^2 \theta}{2\mu_0 k_B \cdot T} \right]$$







## 4 Tesla Helmholtz coils on Xmas beam line

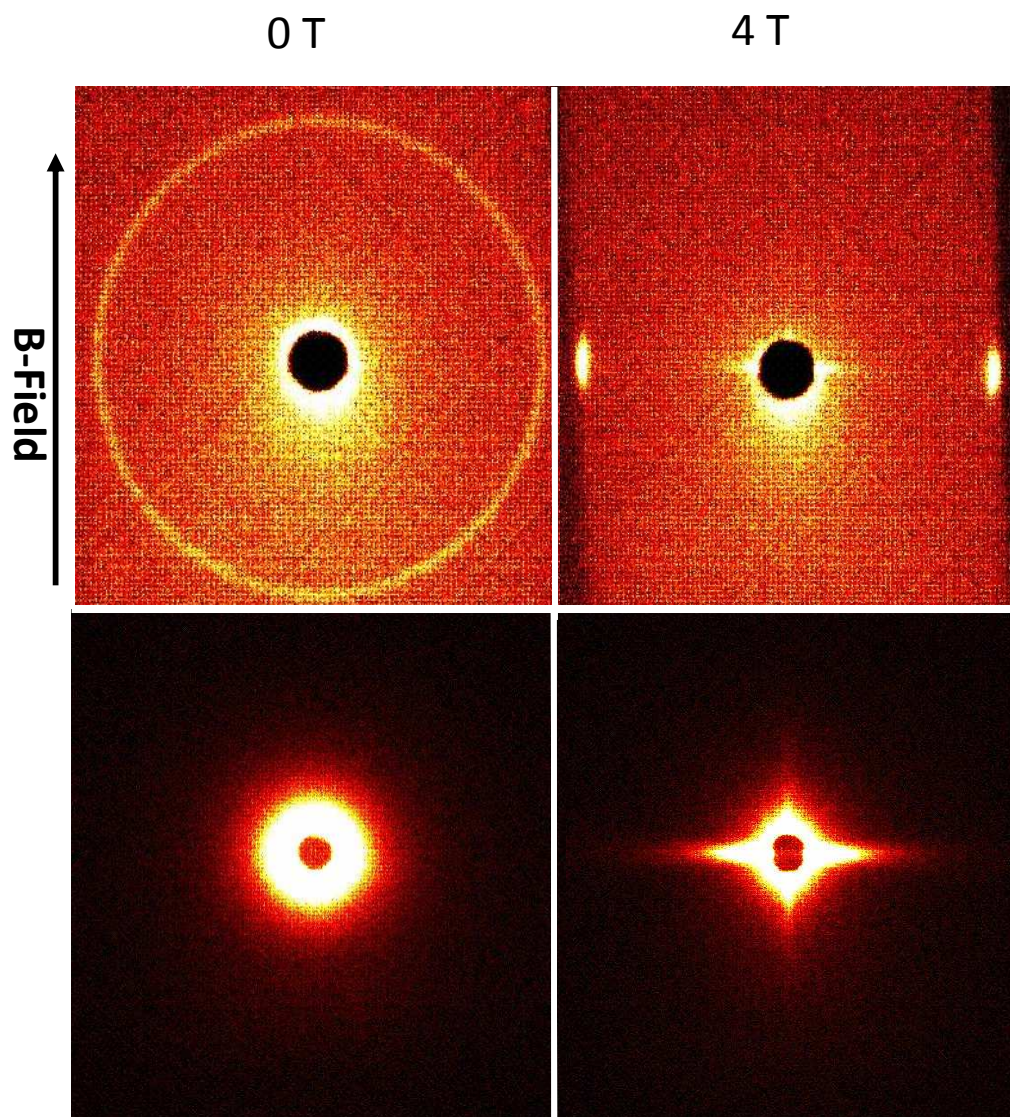


BM28@ESRF





## T6/butanol aggregates in magnetic field ( $B \perp X$ -rays)



S-to-D  $\sim$  1.5m

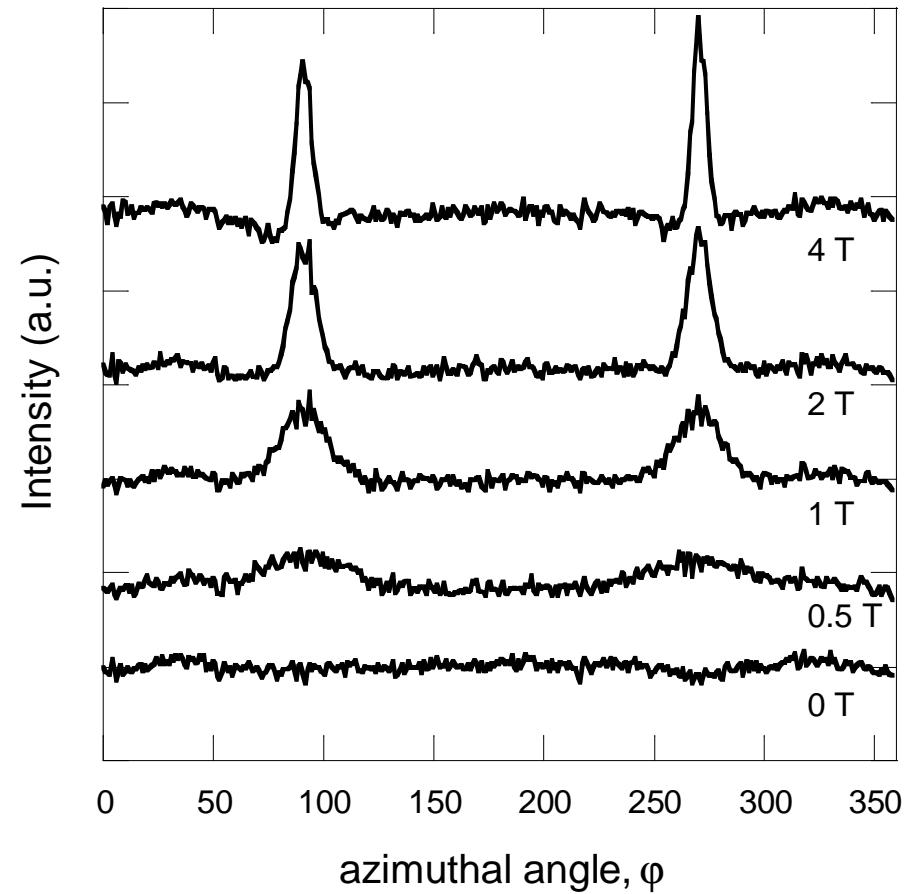
- diffraction ring transforms into diffracted “spots” (arcs)
- molecules “highly” oriented perpendicular to the field

S-to-D  $\sim$  8m

- strong anisotropy mainly in the horizontal scattering direction (perpendicular to the field)
- Aspect ratio from Porod  $\sim$  7

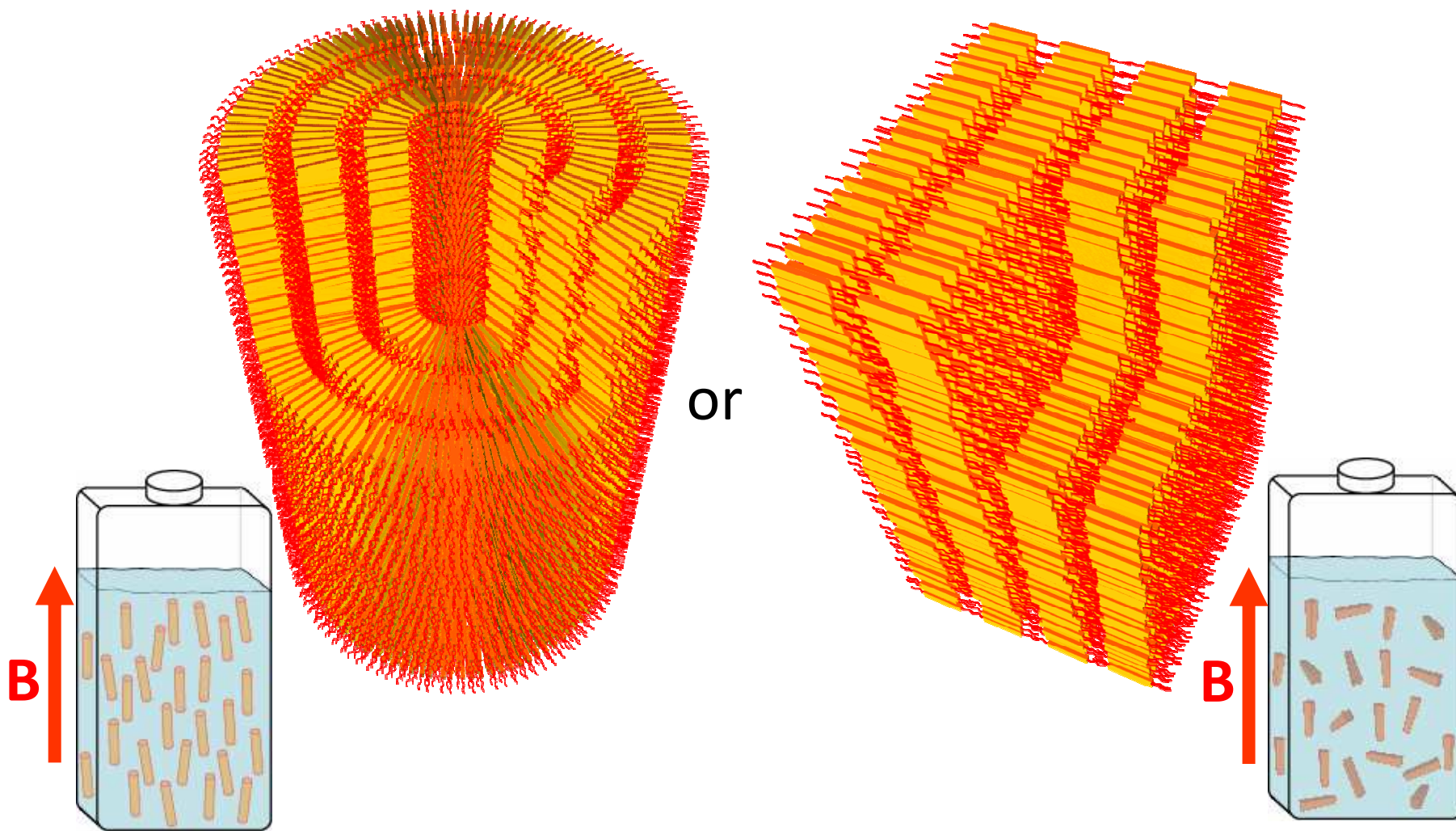


# T6/butanol aggregates in magnetic field ( $B \perp X$ -rays)





# T6/Butanol



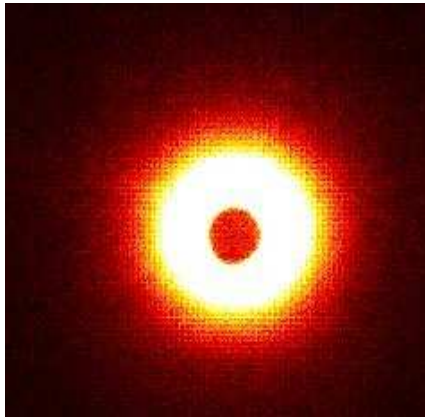
Langmuir 25, 1272 (2009)



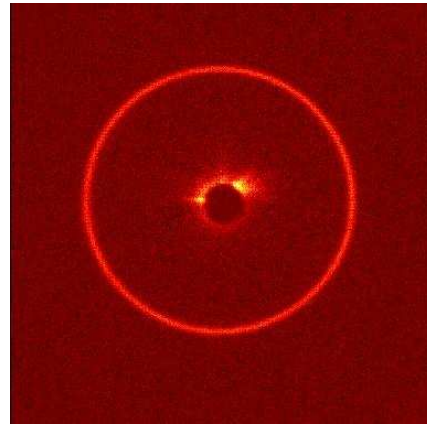
# T6 aggregates in magnetic field ( $B \parallel$ X-rays)

2T

$B \parallel$  X-rays

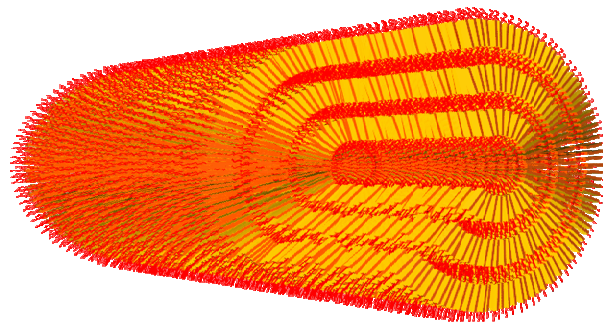


8 m



1.5 m

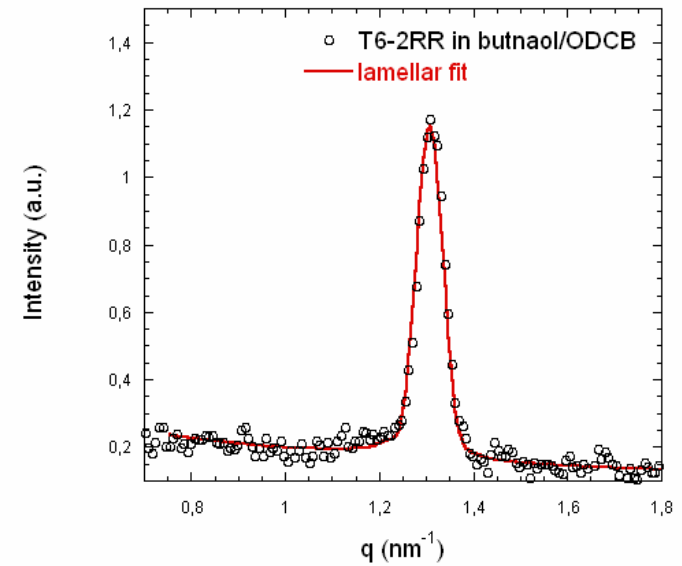
$\leftarrow$  **B**



$\leftarrow$  X-rays

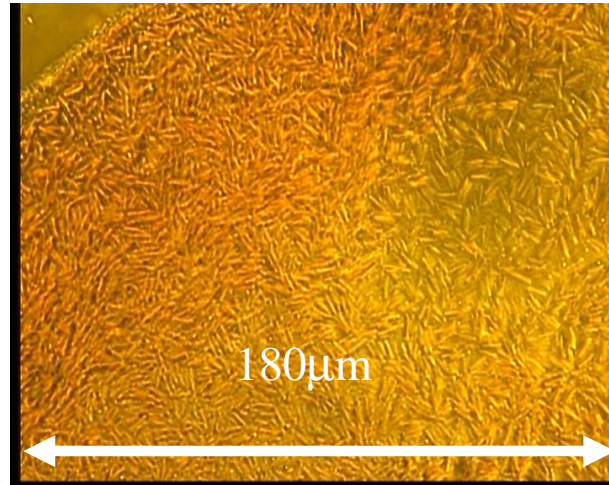
Debye-Scherrer:

$$\Lambda = 72.4 \text{ nm} \rightarrow 15 \text{ T6 layers}$$

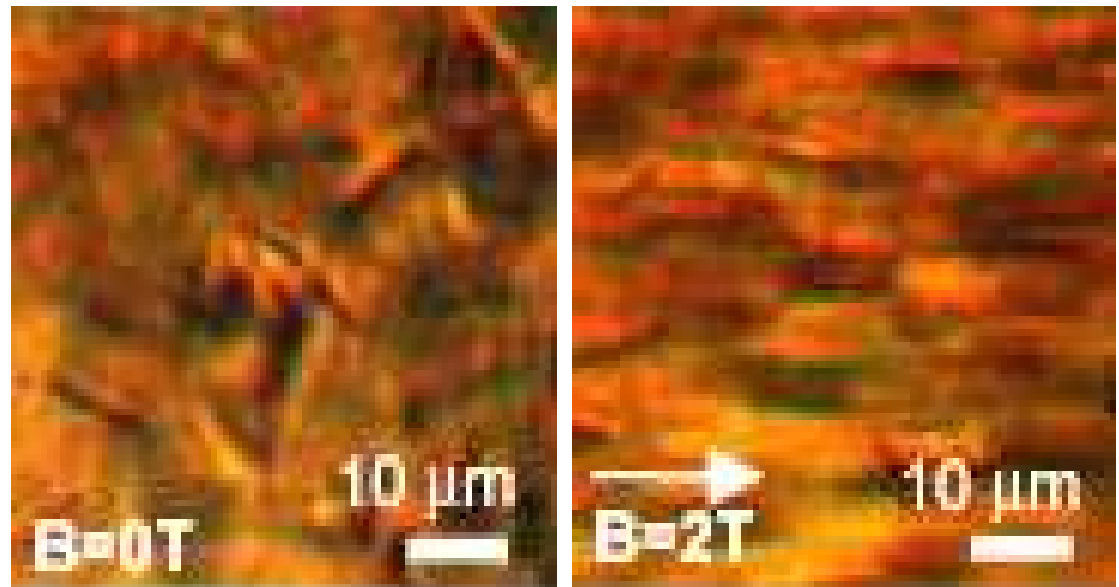


**15.3 lamellar layers**  
**( $l_1 = 3.2 \text{ nm}$ ,  $l_2 = 1.6 \text{ nm}$ )**

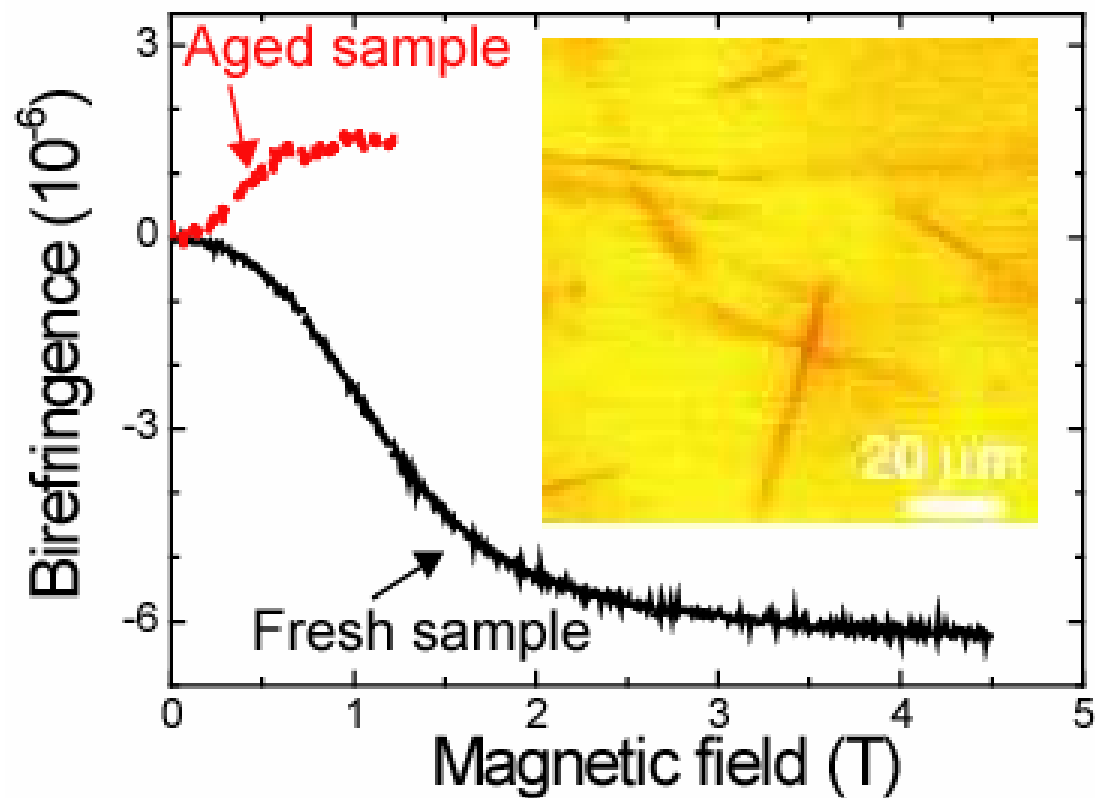
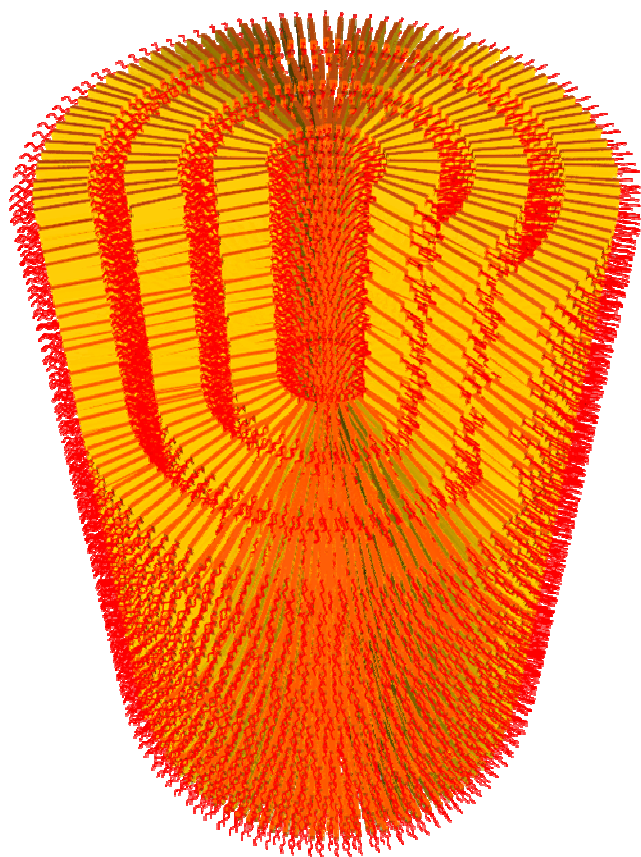
Optical confirmation (sample aged for 3 months)



In-situ confocal microscopy



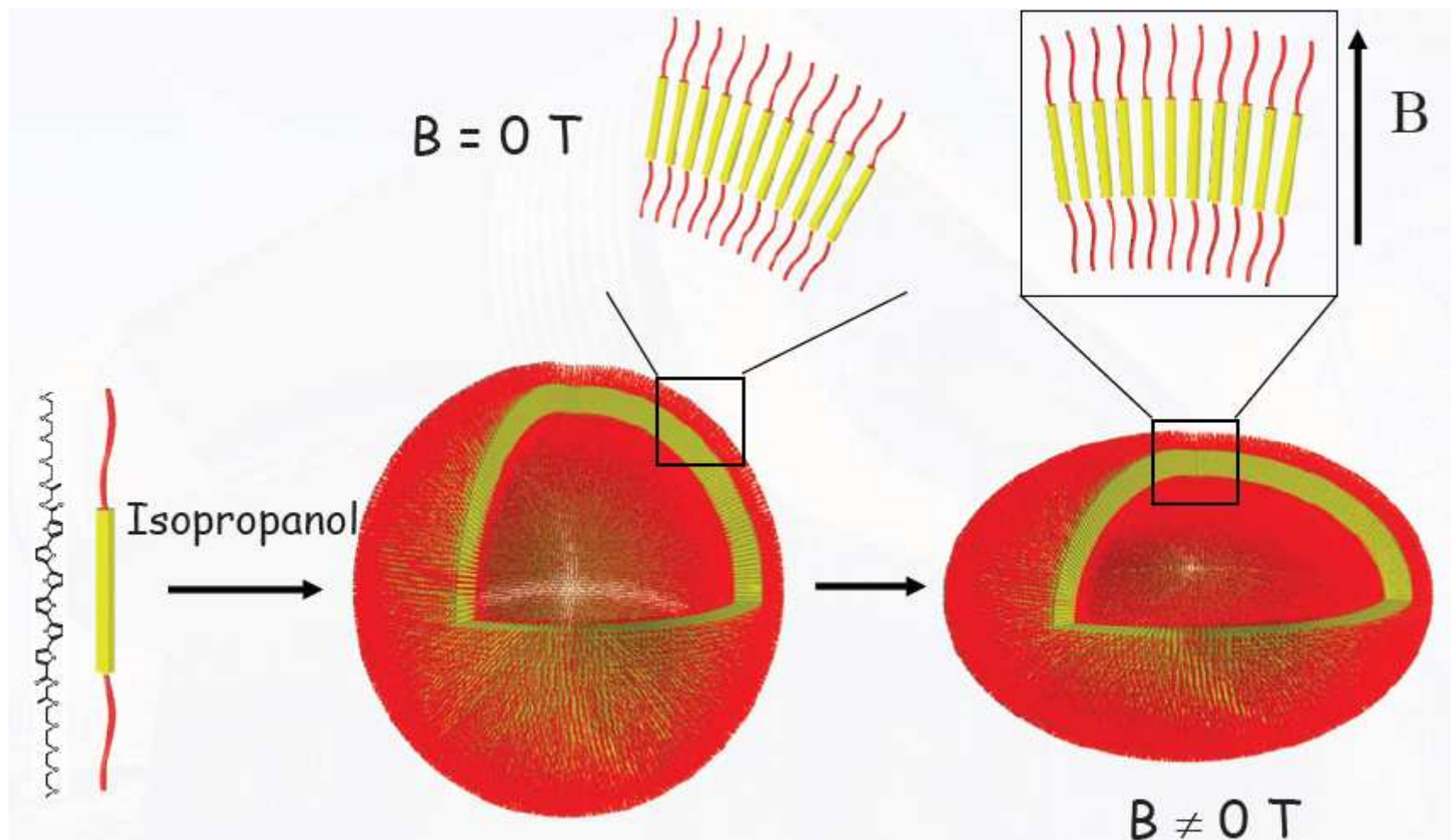
# Cylindrical aggregates



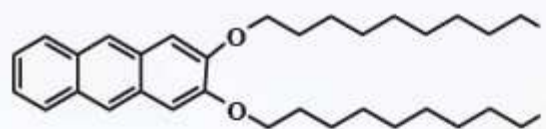




# Magnetic deformation of T6 nanocapsules



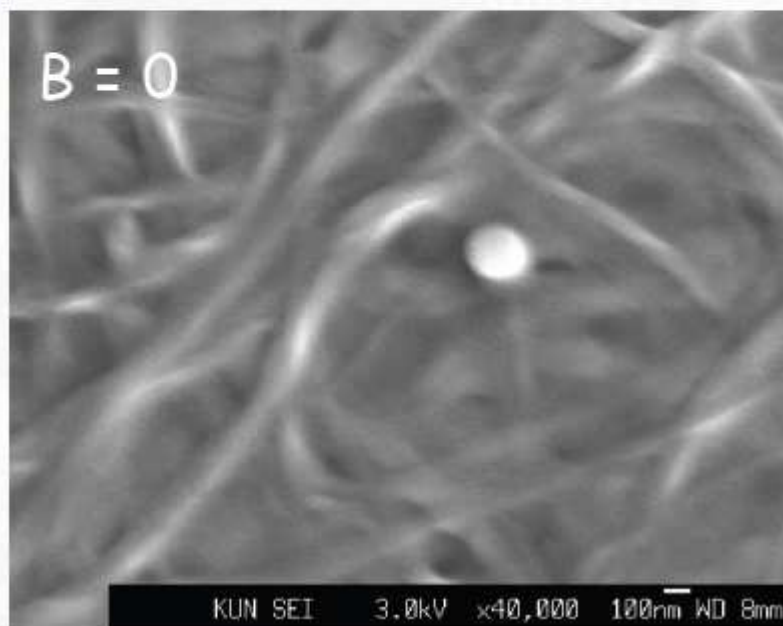


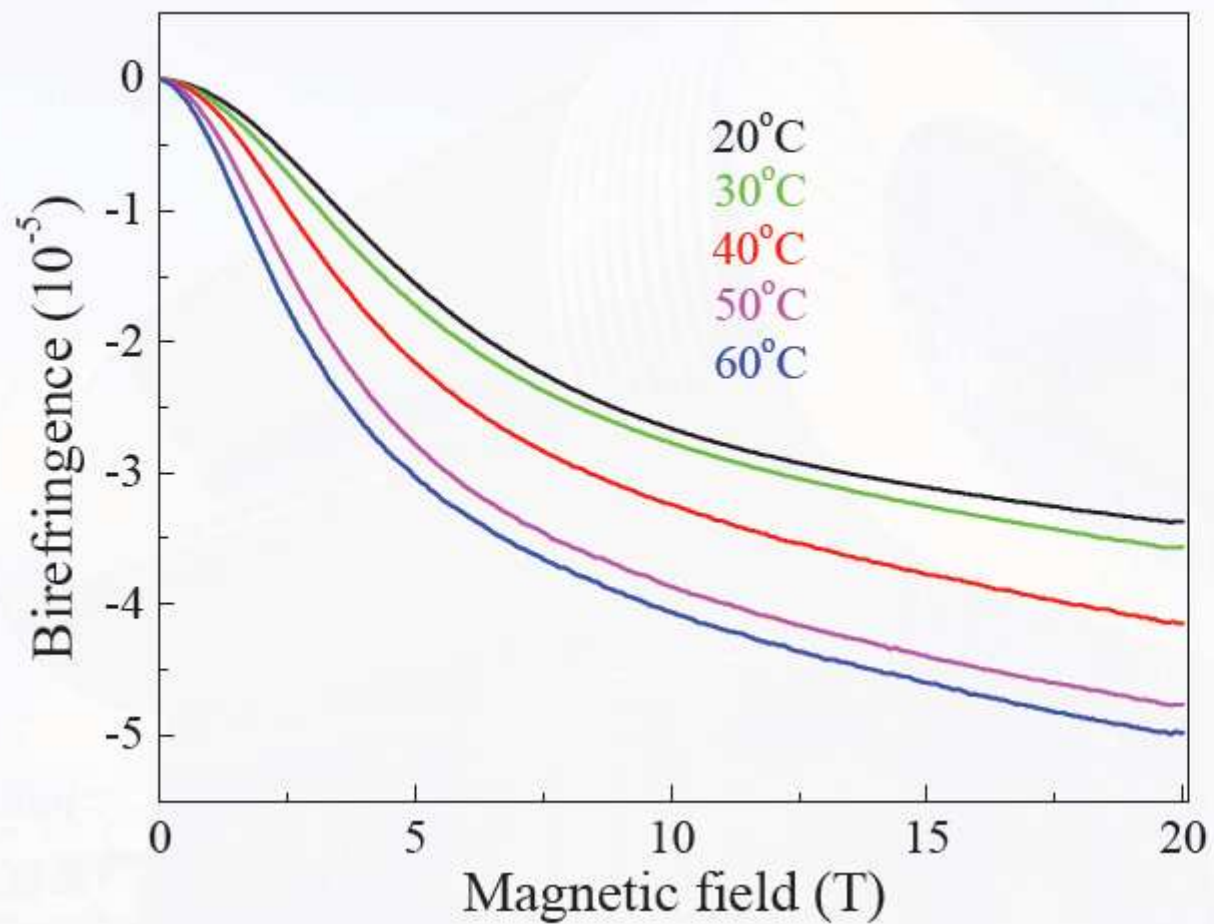


### Gelators

- of alcohols
- forming fiber network

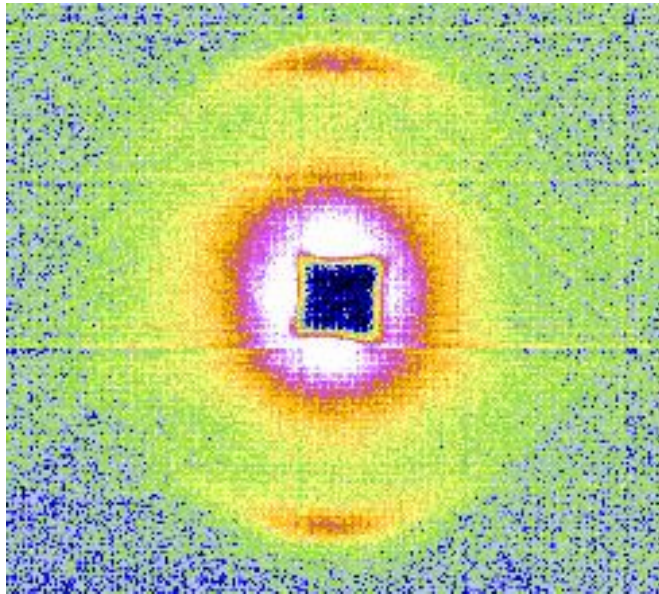
J-P. Desvergne et al., JACS 6, 416 (1991)



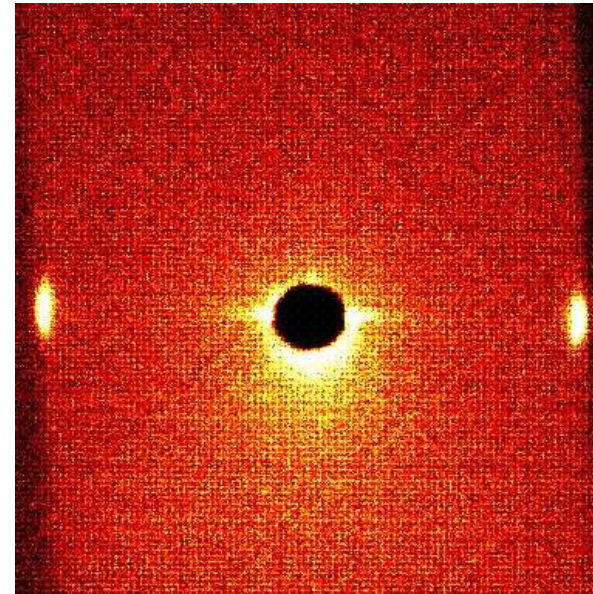


raw birefringence  
data  $\Delta n$

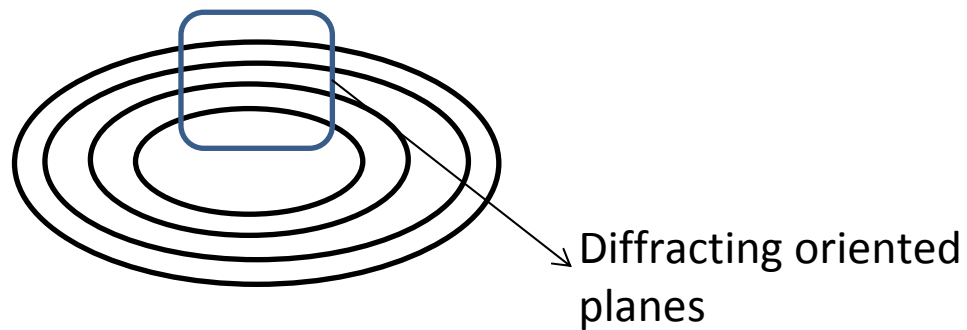
Deformed T6/isopropanol vesicles



Aligned T6/Butanol cylinders



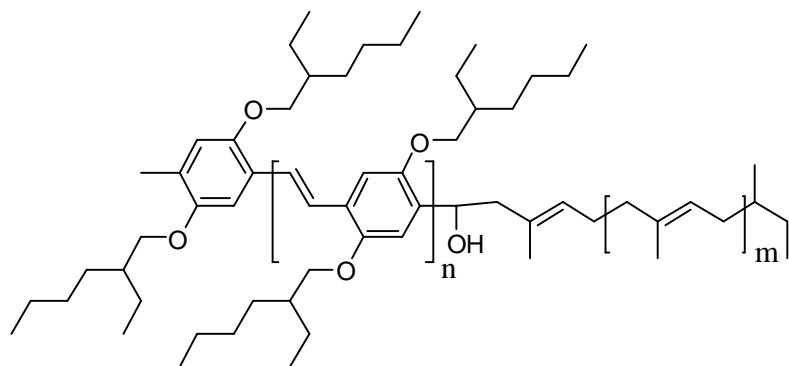
B-Field ↑





# PPV-PI rod-coil block copolymer in 7T magnet

Poly(2,5-di(2'-ethylhexyloxy)-1,4-phenylene vinylene)-*b*-polyisoprene



**Rod - Coil**

$T_{ODS} = 115 \text{ }^{\circ}\text{C}$   
Smectic phase above  $60 \text{ }^{\circ}\text{C}$



7T magnet @ DUBBLE

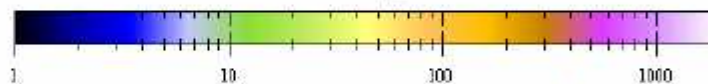
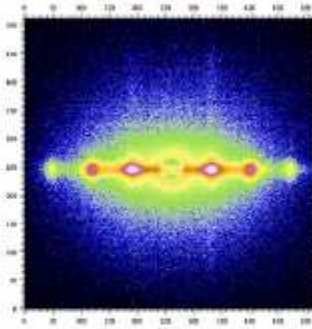
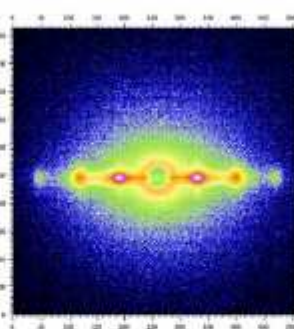
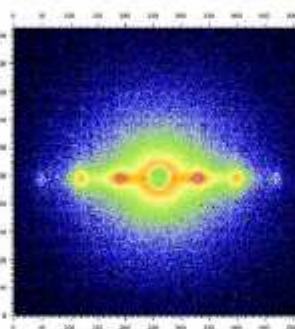
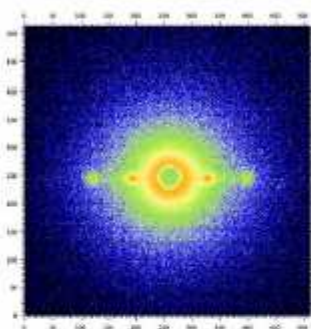
## Quench 150 C => 105 C

138C – 1 min

129C – 2 min

123C – 3 min

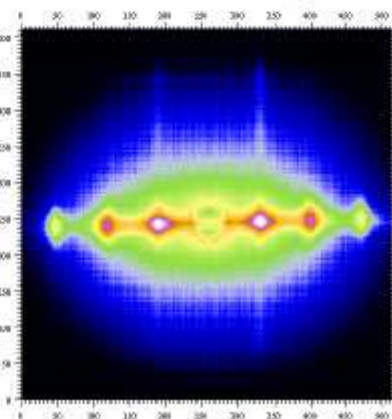
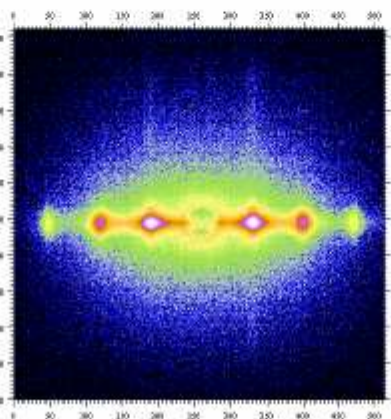
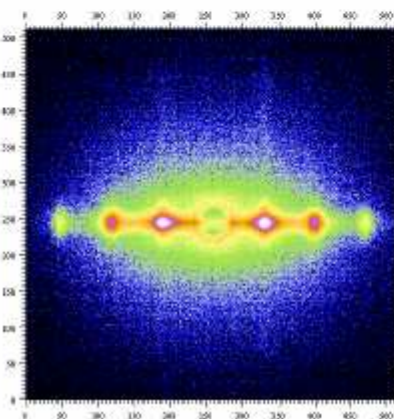
108C – 7 min



105C – 14 min

104C – 21 min

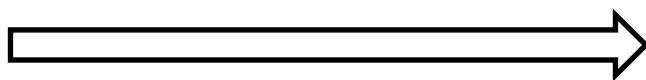
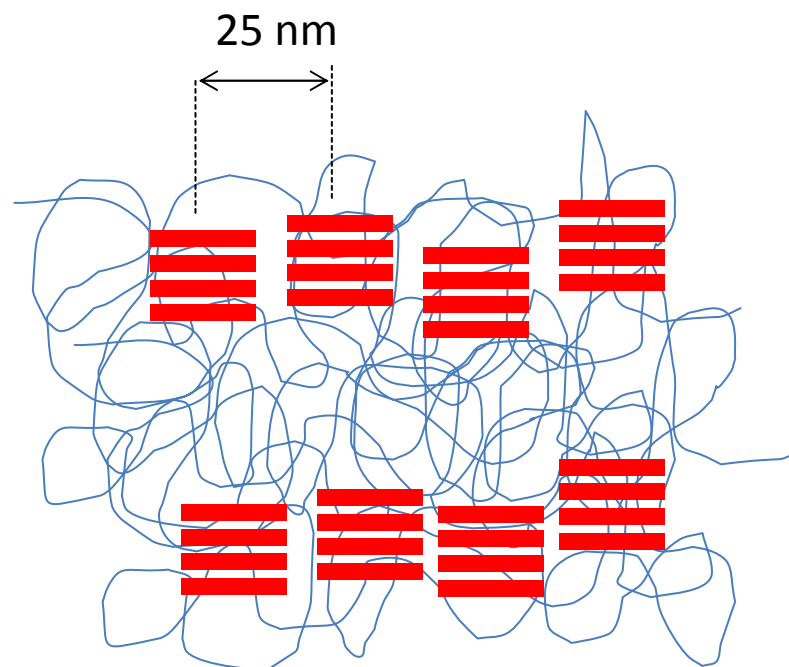
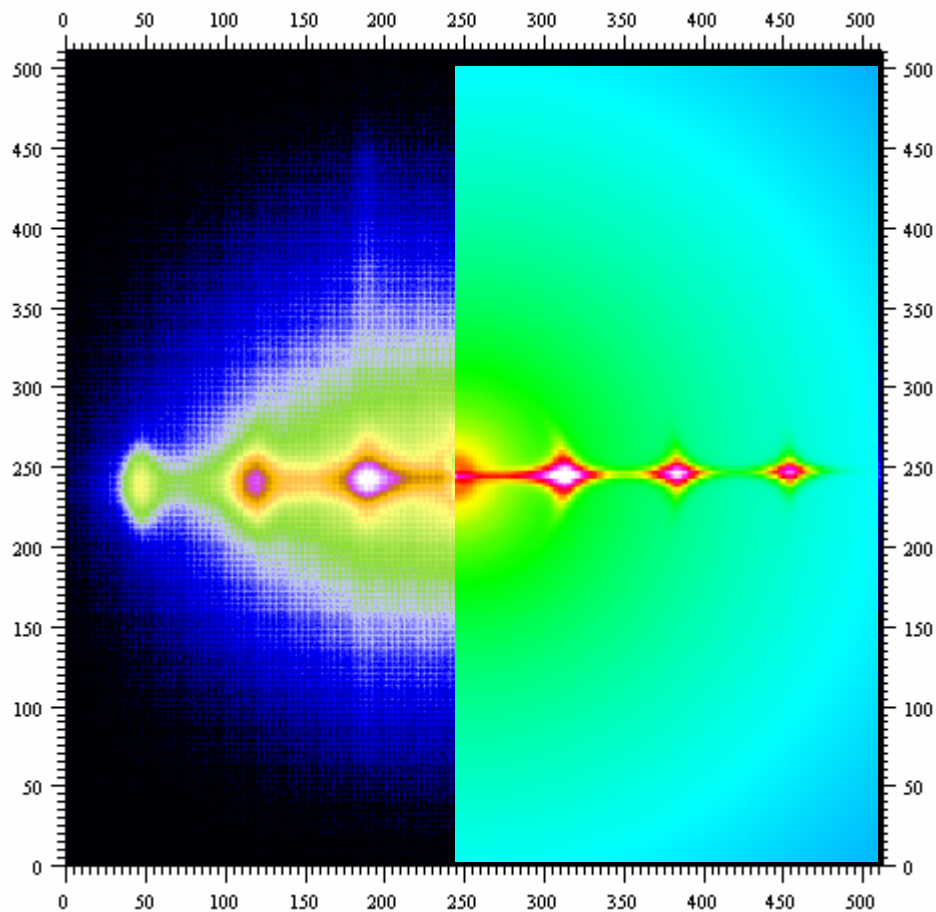
Intensity 107C – 28 min (10 min exposure)



D21-24



# PPV-PI rod-coil block copolymer in 7T magnet

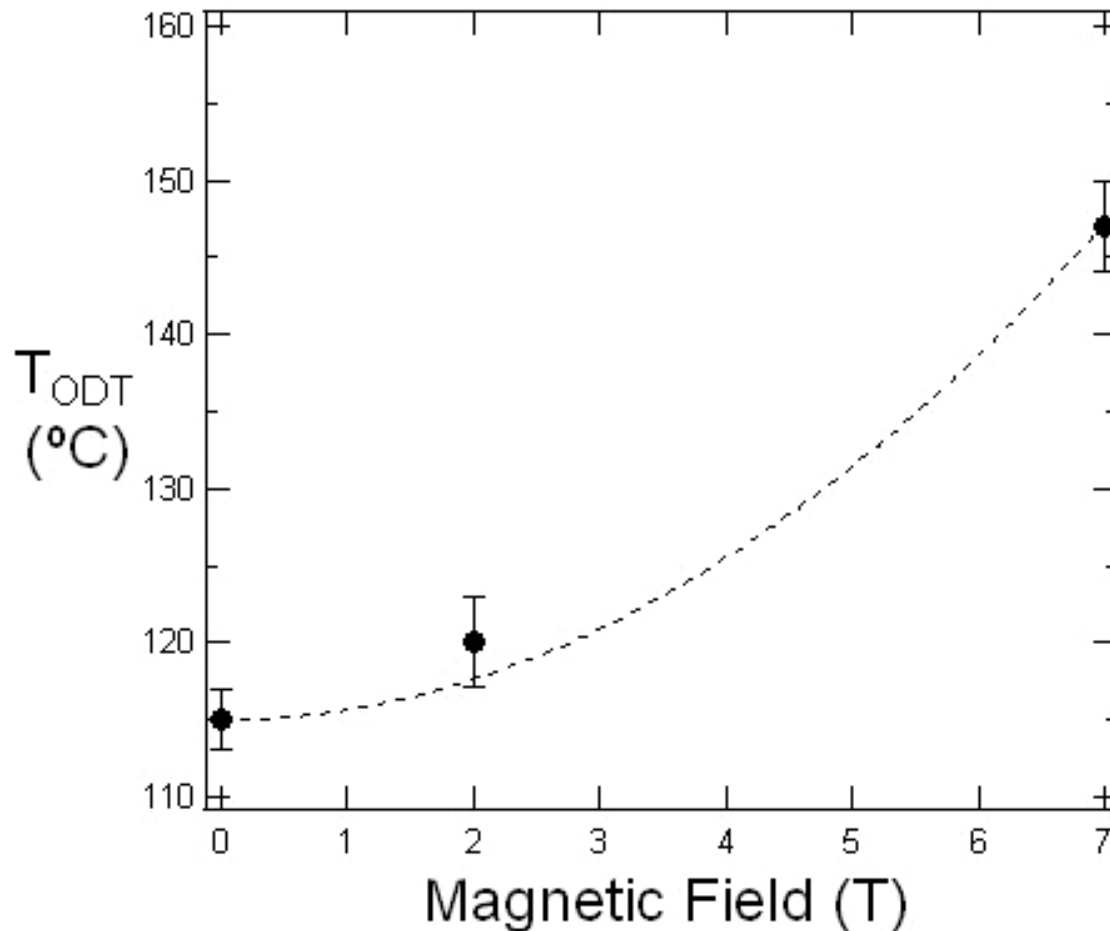


$B = 7T$



## PPV-PI rod-coil block copolymer in 7T magnet

### $T_{ODT}$ influenced by magnetic field



Application of magnetic field stabilizes the smectic phase, extending its T-range of existence



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# Conclusions

- Manipulation of soft-matter coupled with SAXS/WAXS helps in understanding complex structures
- Could be used to control orientation in devices from advanced functional materials ( $\pi$ -cojugated structures)
- Pulsed (high) magnetic fields coupled with XFELs may help solving structure of single molecules (i.e. proteins, small aggregates etc.)





# Acknowledgments



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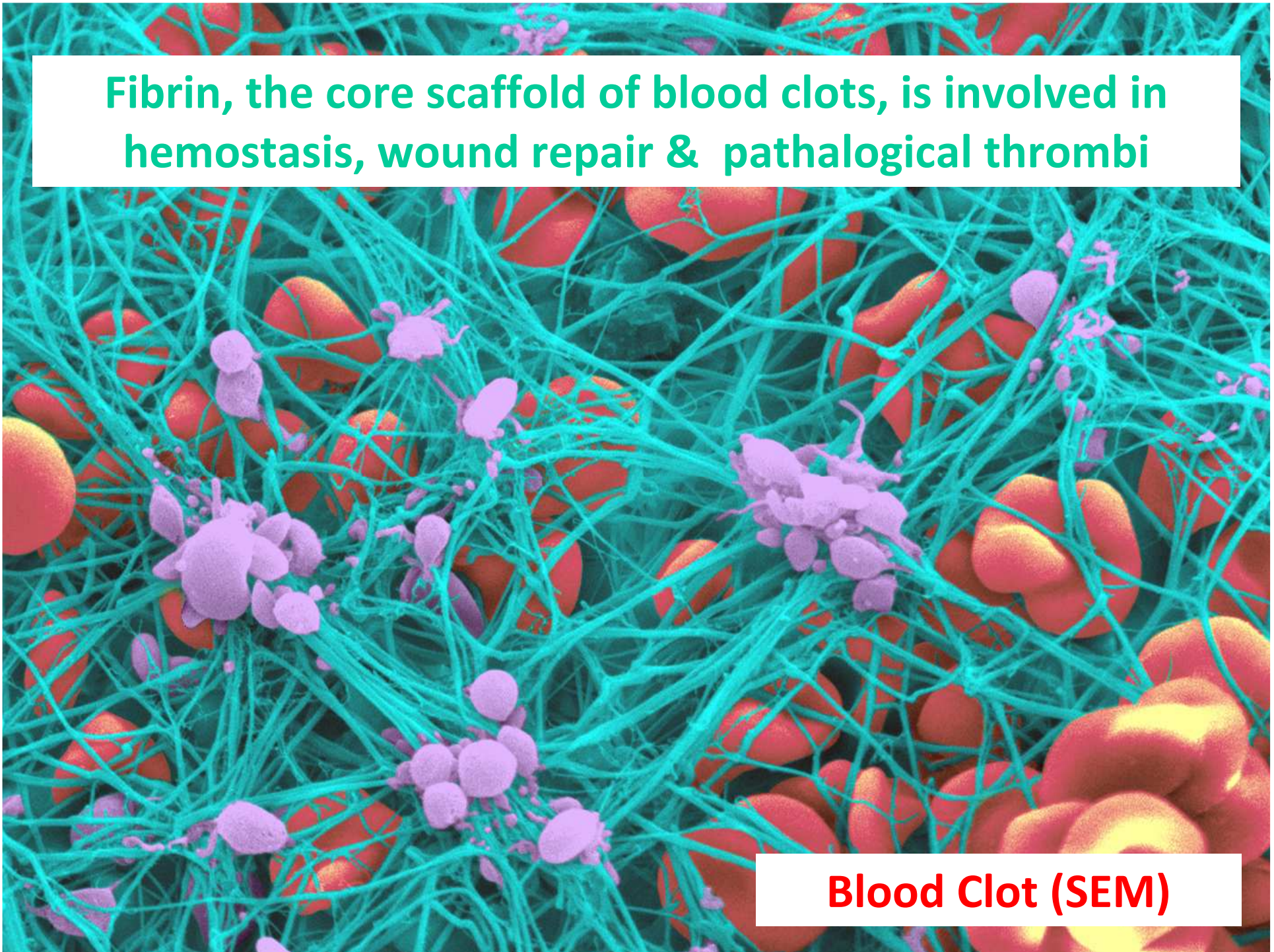
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Laurence Bouchenoire



Bert Meijer  
Albert Schenning  
Martin Wolffs

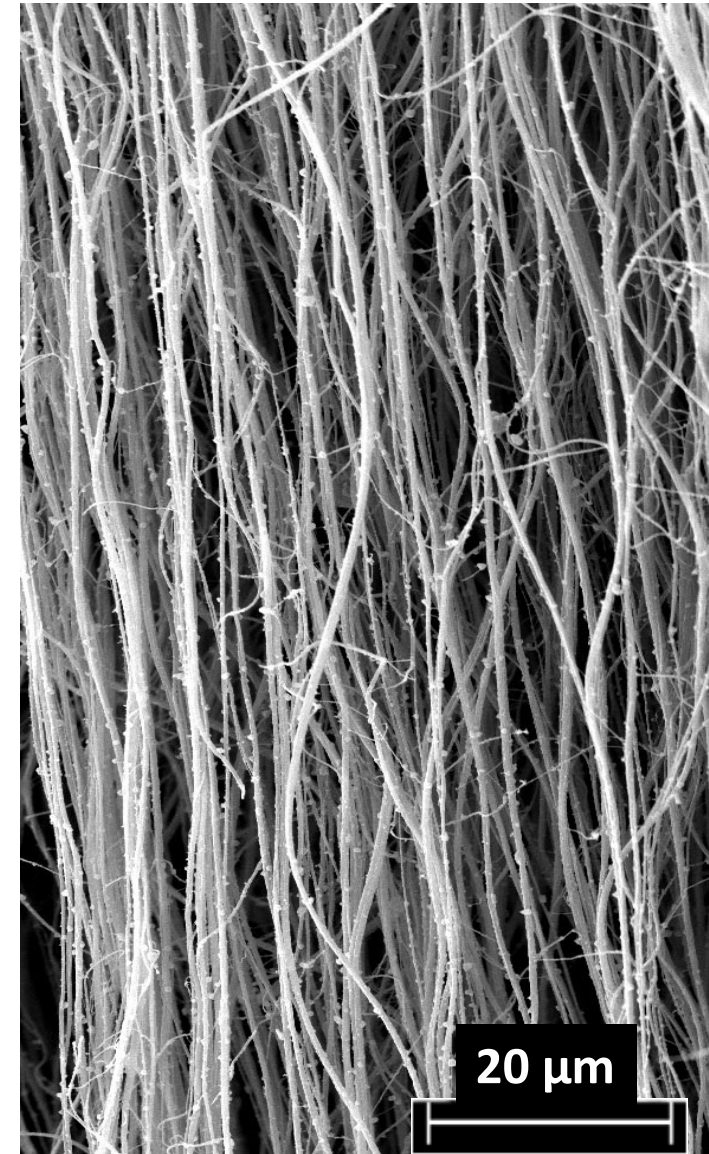
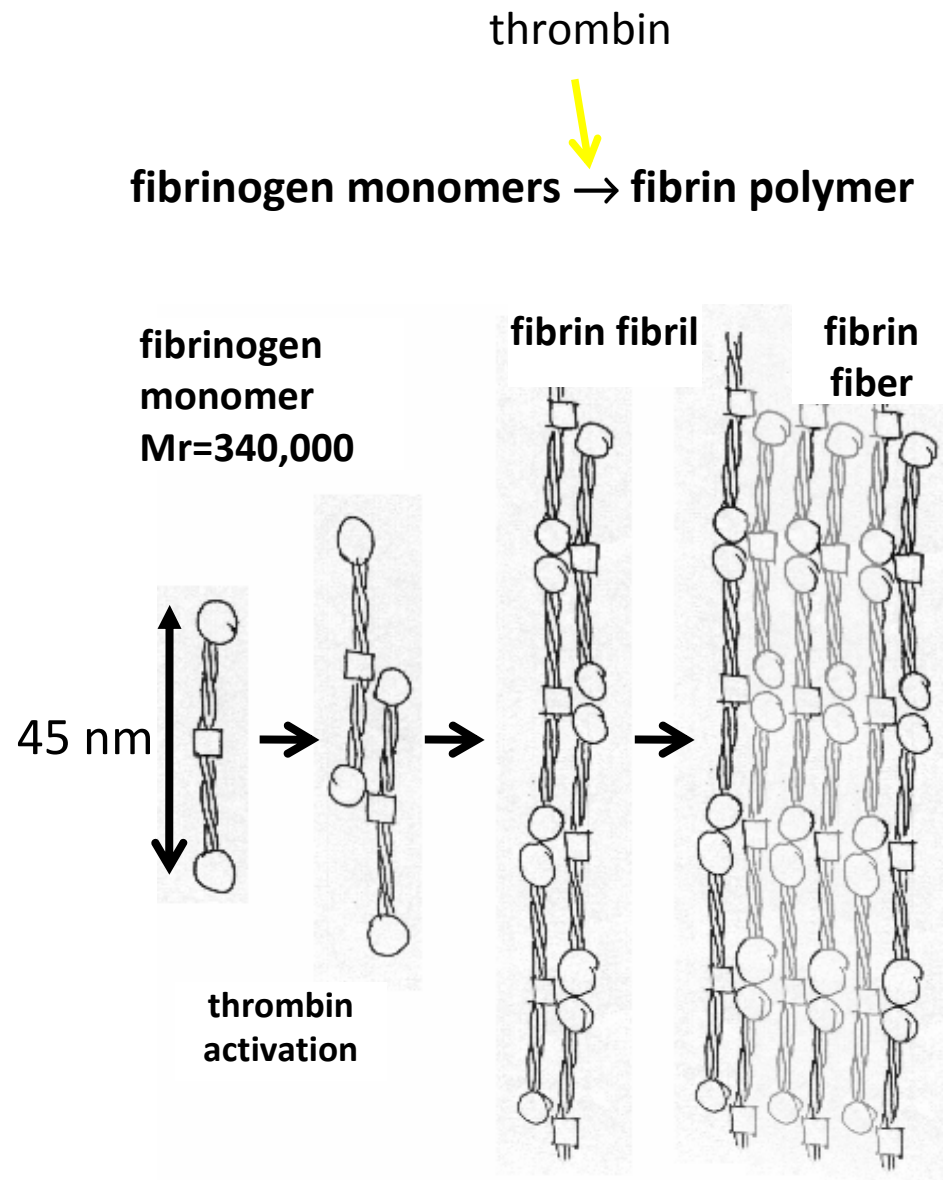
**Fibrin, the core scaffold of blood clots, is involved in hemostasis, wound repair & pathological thrombi**



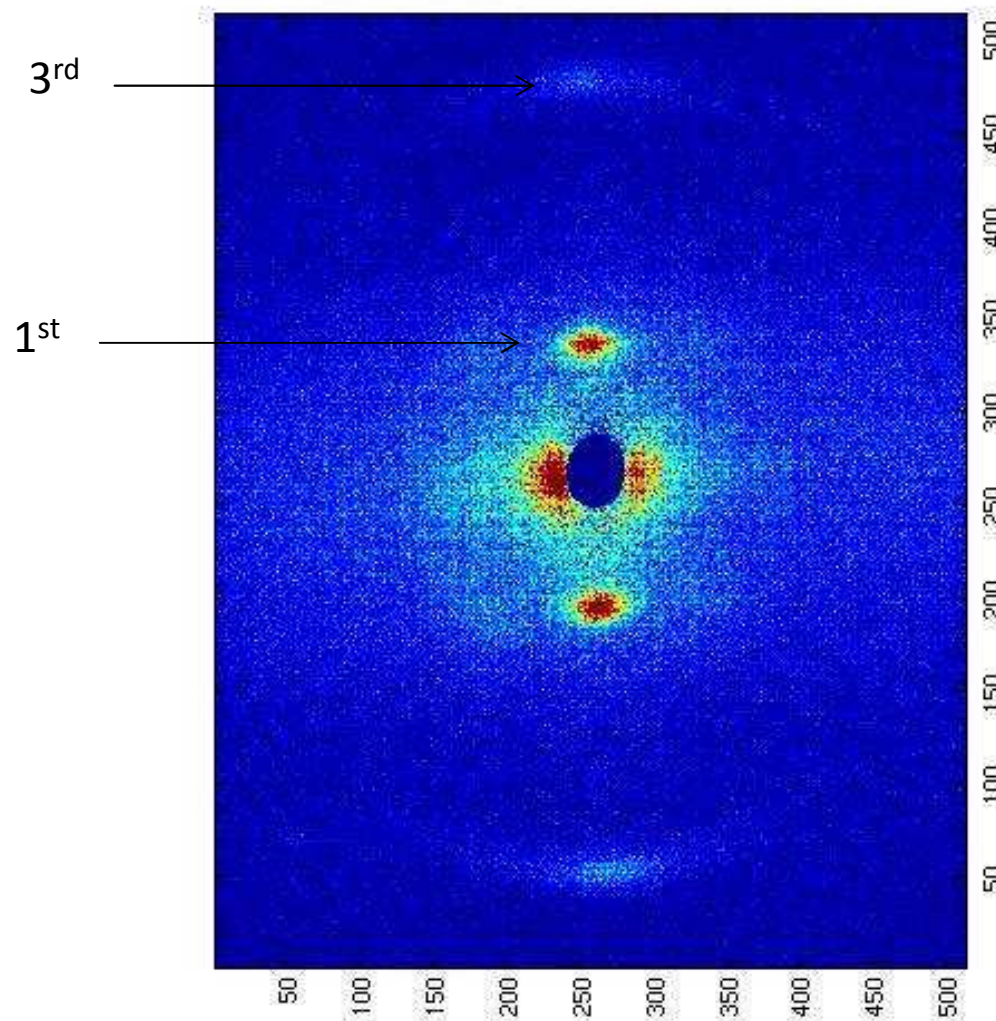
**Blood Clot (SEM)**



# Human blood plasma clot assembled in 7 T



## X-ray diffraction of Fibrin assembled in 12T magnet

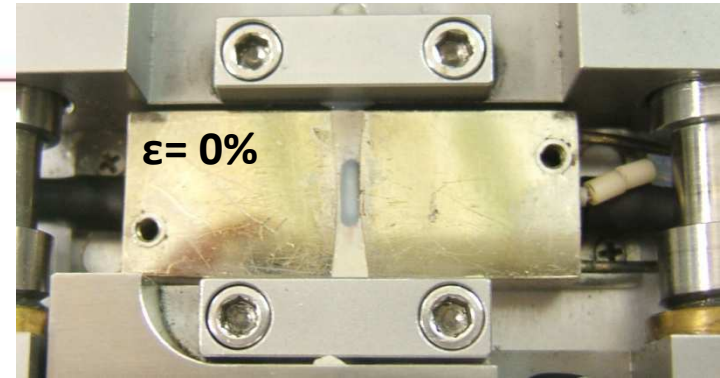


$$d = 22\text{nm}$$

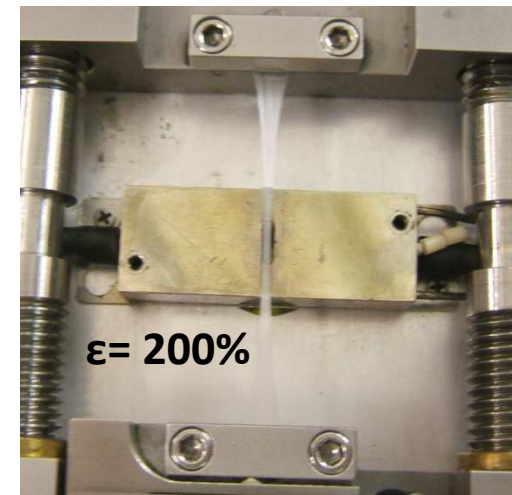


Why magnetic orientation: The near perfectly aligned gels means that all fibers are simultaneously stretched to the same degree.

Why stretch: Fibrin is subjected to different degrees of tensile stress: in the short term, by platelet induced contraction during clot consolidation, and cell traction, during wound repair, &, in the long term, by the shear caused by blood flow round pathological thrombi.

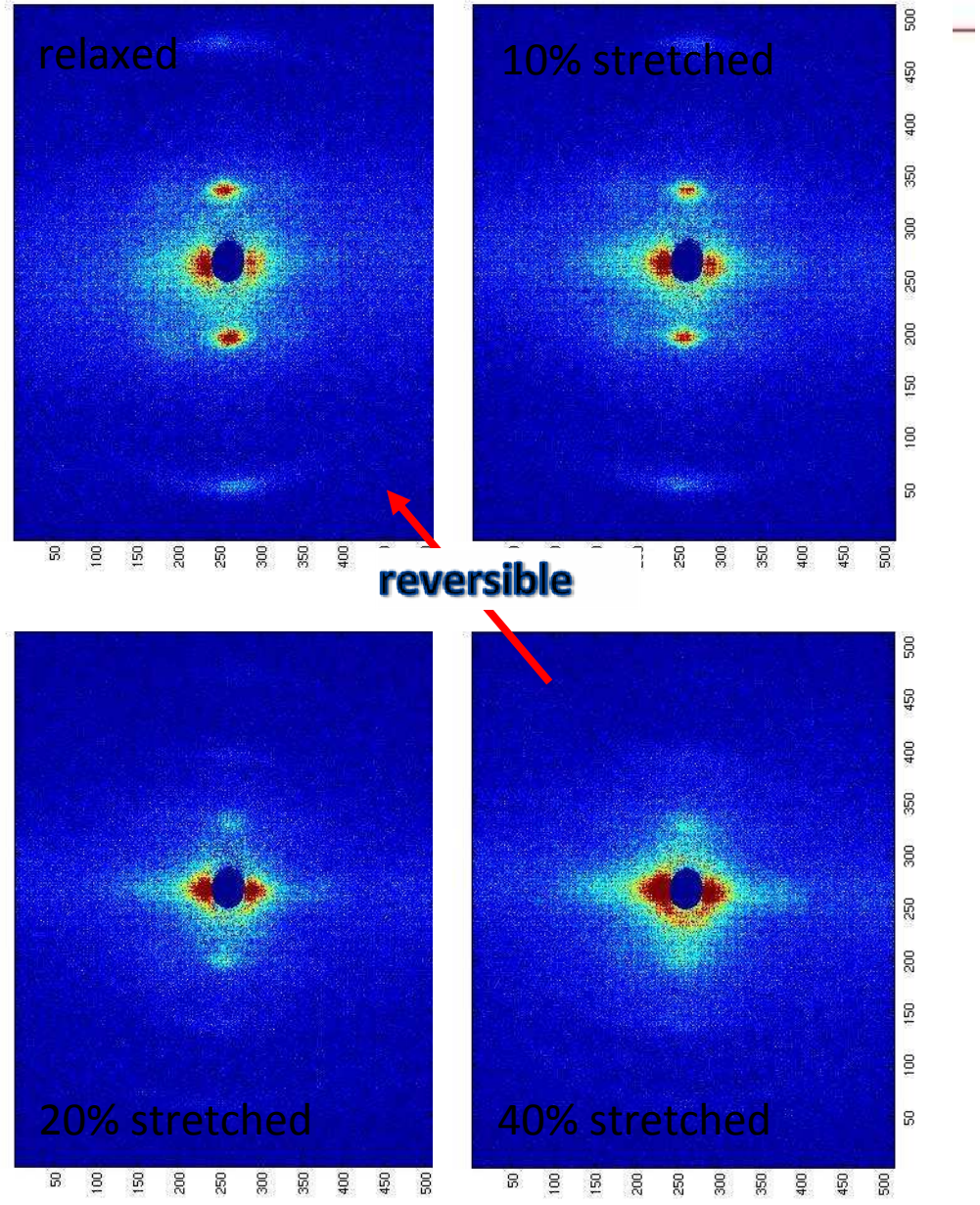
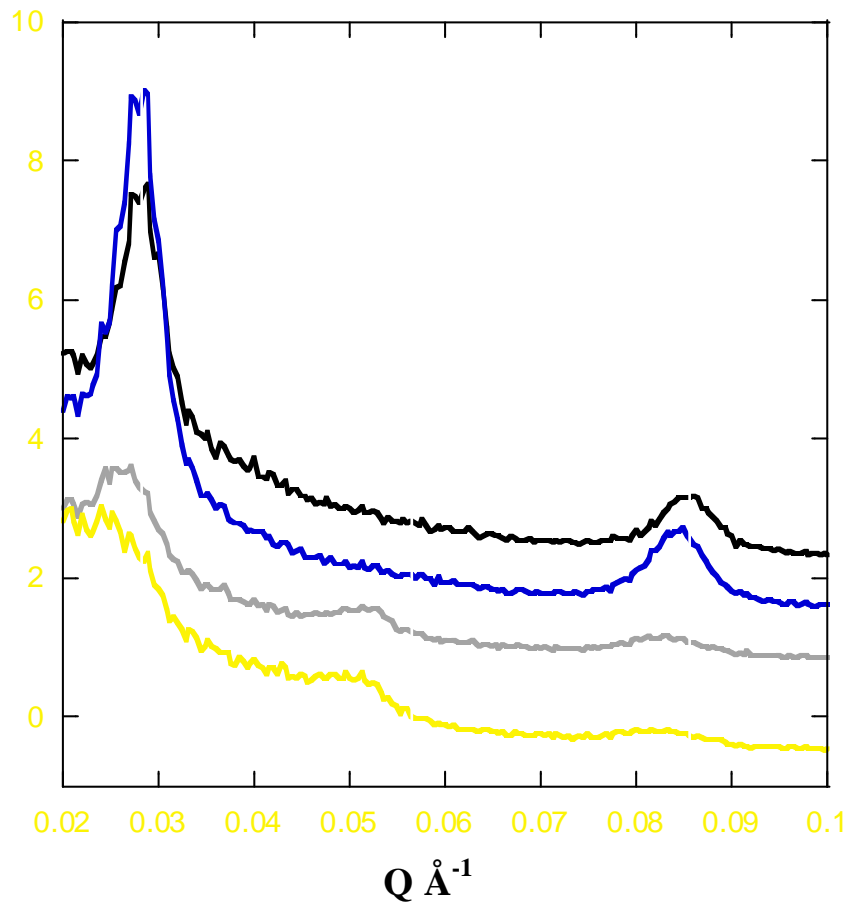


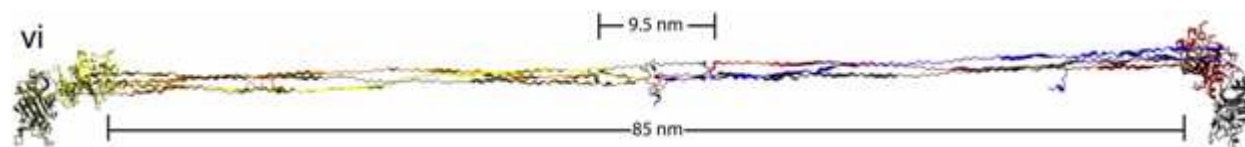
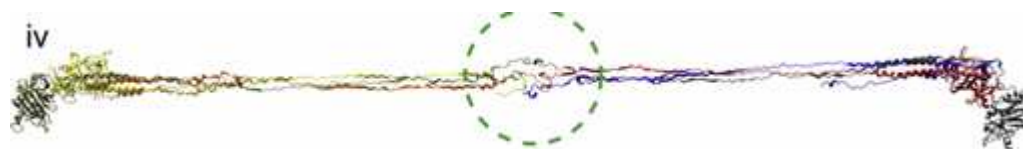
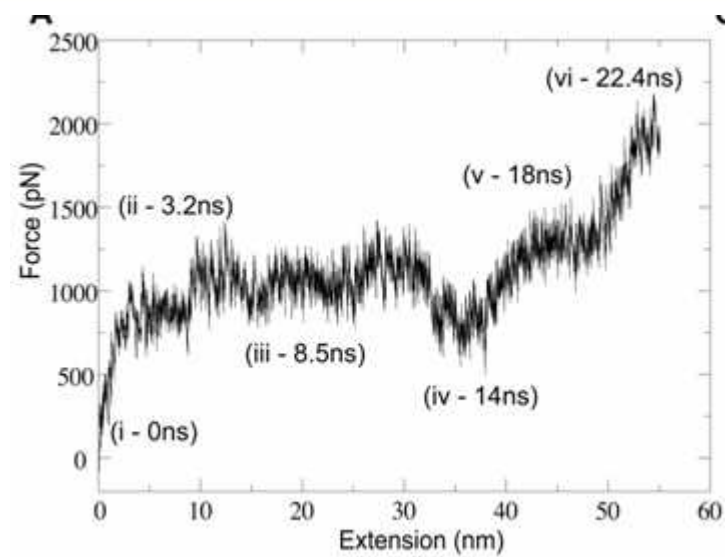
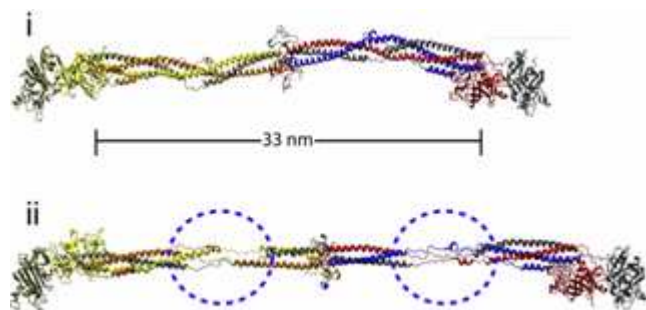
Human platelet free plasma  
clot assembled in 12 T



Stretching device in a water  
saturated atmosphere & X-  
rayed (BM26, DUBBLE)

# Forced unfolding of fibrin







- The initial half-stagger axial repeat is only stretched by about 5% when it disappears at an extension,  $\epsilon \approx 40\%$ .
- A new Bragg peak appears  $\epsilon \approx 20\%$  belonging to a different lattice, indicating the formation of a new ordered transition state which lingers beyond  $\epsilon = 40\%$ .
- On relaxation clots stretched beyond 40% recover to their original length giving axial diffraction again but, surprisingly, the repeat is shorter than the initial value. The structure of stretched relaxed samples is different and the difference increases with the number of stretch-relaxation cycles and the maximum value of  $\epsilon$  experienced. After stretching to  $\epsilon=200\%$  a Bragg spacing of only 18nm was recorded.